

# Paper Review of Data-assisted reduced-order modeling of extreme events in complex dynamical systems

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# Problem Setup

- $\frac{du}{dt} = F(u) = Lu + h(u)$ ,  $u(t) \in \mathbb{R}^n$ , where  $n$  is a large number.
- We are specifically interested in systems whose dynamics results in a intrinsically low dimensional globally attracting manifold.

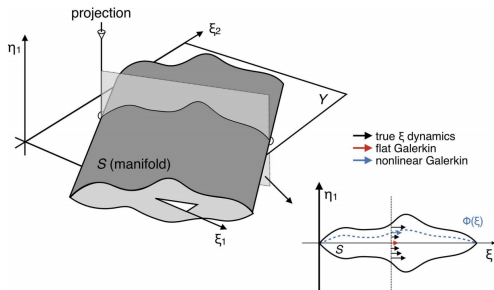


Figure: Example of a Globally Attracting Manifold

# Problem Setup

Goal:

- Dimensionality Reduction of the ODE
- Derive an ODE in a low dimensional subspace, which preserves as much dynamical property of original ODE as possible.

Ansatz:

$$u = \mathbf{Y}\xi + \mathbf{Z}\eta + b \quad (1)$$

where the columns of matrix  $\mathbf{Y} = [y_1, \dots, y_m]$  form an orthonormal basis of  $Y$ , an  $m$ -dimensional subspace of  $\mathbb{R}^d$ , and the columns of  $\mathbf{Z} = [z_1, \dots, z_{d-m}]$  make up an orthonormal basis for the orthogonal complement  $Z = \mathbb{R}^d \setminus Y$ ;  $\xi$  and  $\eta$  are the projection coordinates associated with  $\mathbf{Y}$  and  $\mathbf{Z}$ .

# Problem Setup

How to get  $\mathbf{Y}$  and  $\mathbf{Z}$ ?

- Randomly sample 10000 points on the attractor, denoted as  $\mathbf{X}$ .
- Do PCA on these datapoints:  
Let  $C = X^T X$ , then  $\mathbf{Y} = [y_1, \dots, y_m]$ , where  $[y_1, \dots, y_m]$  are the eigenvectors of  $C$  corresponding to the  $m$  largest eigenvalues of  $C$ .  
 $\mathbf{Z} = [z_1, \dots, z_{d-m}]$ , where  $[z_1, \dots, z_{d-m}]$  are the rest of eigenvectors of  $C$ .

Back to equation (1), we know  $Y$ ,  $Z$  and the ODE for  $u$ , can we have the ODE for  $\xi$ ?

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Plug (1) back to the ODE, we get

$$\frac{d\xi}{dt} = Y^T L Y \xi + Y^T L Z \eta + Y^T h(Y \xi + Z \eta + b) + Y^T L b \quad (2)$$

If  $|\eta| \ll |\xi|$ , then we may assume  $\eta = 0$ , leading to a  $m$ -dimensional system:

$$\frac{d\xi}{dt} = Y^T L Y \xi + Y^T h(Y \xi + b) + Y^T L b = F_\xi(\xi) \quad (3)$$

This is called the Flat Galerkin Method.

# Nonlinear Galerkin Projection

Problems:

- $\eta = 0$  may be too strict.
- $Z$  is derived merely based on statistical properties of the manifold without addressing the dynamics. This implies that even if  $\eta$  has small magnitude on average it may play a big role in the dynamics of the space.

Existing solution: Let  $\eta = \Phi(\xi)$ , then

$$\frac{d\xi}{dt} = Y^T L Y \xi + Y^T L Z \Phi(\xi) + Y^T h(Y \xi + Z \Phi(\xi) + b) + Y^T L b \quad (4)$$

This is called the Nonlinear Galerkin Projection Method.



# Nonlinear Galerkin Projection

Problems:

- $\Phi$  has to be found empirically.
- $\Phi$  may not be well-defined!

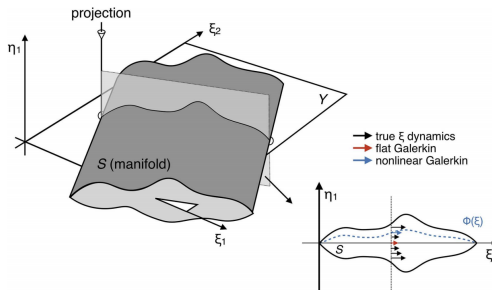


Figure: Example of a Globally Attracting Manifold

Proposed Data-driven solution:

$$\frac{d\xi}{dt} = Y^T L Y \xi + Y^T L Z \eta + Y^T h(Y \xi + Z \eta + b) + Y^T L b \quad (5)$$

rewritten as

$$\frac{d\xi}{dt} = F_\xi(\xi) + G(\xi, \eta) \quad (6)$$

Try to learn  $G$  from the data.

Assumption:  $\Psi(t) = G(\xi(t), \eta(t)) \approx \hat{G}(\xi(t), \xi(t - \tau), \xi(t - 2\tau), \dots)$

What is  $\hat{G}$ ? LSTM...

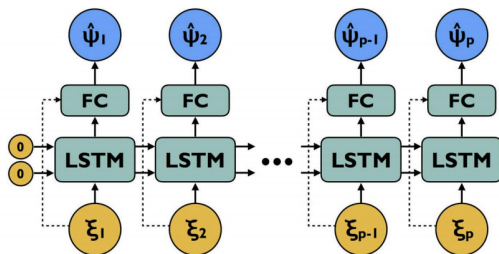


Figure: Architecture of the First Proposed Model

$$\text{Loss: } L = \sum_{i=1}^p w_i \|\hat{\Psi}_i - \Psi_i\|^2, \text{ where } w_i = \begin{cases} w_0 & 0 < i \leq p_t \\ 1 & p_t < i \leq p \end{cases}$$

# LSTM

Problem:

- Input representing the  $\xi$  is always accurate regardless of any errors made in predicting the dynamics previously, i.e. the model only learns to predict one step ahead.
- This is undesirable especially for chaotic systems where errors tend to grow exponentially.

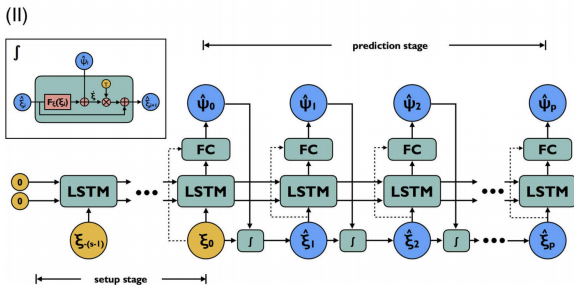


Figure: Architecture of the Second Proposed Model

- Loss for the second model:  
$$L = \sum_{i=1}^P w_i \|\hat{\Psi}_i + F_{\xi}(\hat{\xi}_i) - \xi_i\|^2, \text{ where } w_i = \gamma^{i-1}, 0 < \gamma < 1.$$
- Use first model for pre-training, use second model for fine-tuning and prediction.

An alternative to the method proposed is fully data-driven modeling, which simply make predictions of  $\xi$  based on previous observed  $\xi$ s. It is shown that this fully data-driven model is worse in the numerical experiments.

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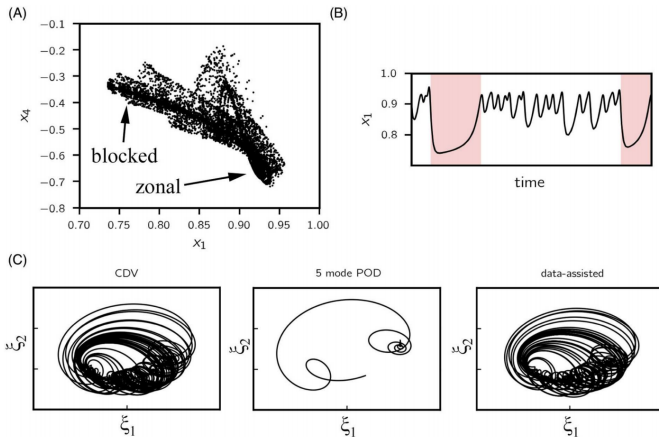
$$\begin{aligned}
 \dot{x}_1 &= \gamma_1^* x_3 - C(x_1 - x_1^*), & \dot{x}_2 &= -(\alpha_1 x_1 - \beta_1) x_3 - C x_2 - \delta_1 x_4 x_6, \\
 \dot{x}_3 &= (\alpha_1 x_1 - \beta_1) x_2 - \gamma_1 x_1 - C x_3 + \delta_1 x_4 x_5, & \dot{x}_4 &= \gamma_2^* x_6 - C(x_4 - x_4^*) + \varepsilon(x_2 x_6 - x_3 x_5), \\
 \dot{x}_5 &= -(\alpha_2 x_1 - \beta_2) x_6 - C x_5 - \delta_2 x_4 x_3, & \dot{x}_6 &= (\alpha_2 x_1 - \beta_2) x_5 - \gamma_2 x_4 - C x_6 + \delta_2 x_4 x_2,
 \end{aligned}$$

where the model coefficients are given by

$$\begin{aligned}
 \alpha_m &= \frac{8\sqrt{2}m^2(b^2 + m^2 - 1)}{\pi(4m^2 - 1)(b^2 + m^2)}, & \beta_m &= \frac{\beta b^2}{b^2 + m^2}, \\
 \delta_m &= \frac{64\sqrt{2}b^2 - m^2 + 1}{15\pi(b^2 + m^2)}, & \gamma_m^* &= \gamma \frac{4\sqrt{2}mb}{\pi(4m^2 - 1)}, \\
 \varepsilon &= \frac{16\sqrt{2}}{5\pi}, & \gamma_m &= \gamma \frac{4\sqrt{2}m^3 b}{\pi(4m^2 - 1)(b^2 + m^2)},
 \end{aligned}$$

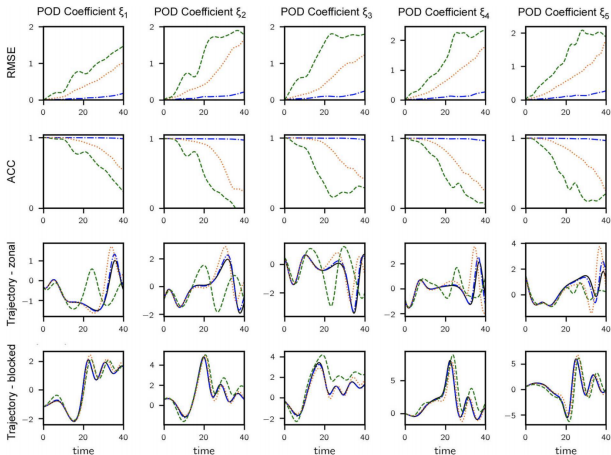
Figure: CDV Equations

$(x_1^*, x_4^*, C, \beta, \gamma, b) = (0.95, -0.760955, 0.1, 1.25, 0.2, 0.5)$ . Use 10000 trajectories, 80% for training, 10% for validation, 10% for test.  $n_{LSTM} = 1$ ,  $n_{FC} = 16$ . Try to predict  $p = 200$  time steps ahead, with  $\delta t = 0.01$ .



**Fig 3. CDV system.** (A)  $10^4$  points sampled from the CDV attractor, projected to  $(x_1, x_4)$  plane. (B) Example time series for  $x_1$ ; *blocked* flow regime is shaded in red. (C) Length-2000 trajectory projected to the first two POD modes (normalized) integrated using the CDV model (left), 5-mode POD projected model (middle) and data-assisted model (right). Despite preserving 99.6% of the total variance, the 5-mode projected model has a single fixed point as opposed to a chaotic attractor. Data-assisted model, however, is able to preserve the geometric features of the original attractor.





**Fig 4. Results for CDV system.** (Row 1) RMSE vs. lead time for 5-mode POD projected model (orange dotted), data-assisted model (blue dash-dotted) and purely data-driven model (green dashed). (Row 2) ACC vs. lead time. (Row 3) A sample trajectory corresponding to zonal flow—true trajectory is shown (black solid). (Row 4) A sample trajectory involving regime transition (happening around  $t = 20$ ). For rows 1, 3 and 4, plotted values are normalized by the standard deviation of each dimension.

# Intermittent bursts of dissipation in Kolmogorov

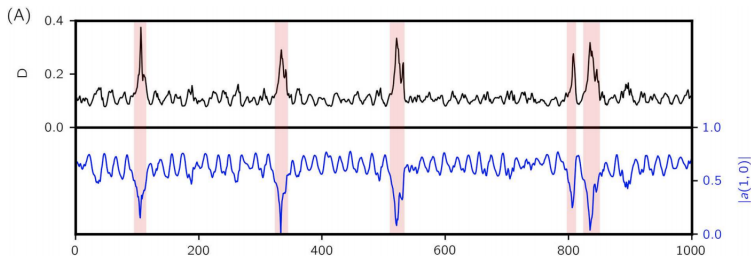
$$\begin{aligned}\partial_t \mathbf{u} &= -\mathbf{u} \cdot \nabla \mathbf{u} - \nabla p + \nu \Delta \mathbf{u} + \mathbf{f} \\ \nabla \cdot \mathbf{u} &= 0\end{aligned}\tag{17}$$

where  $\mathbf{u} = (u_x, u_y)$  is the fluid velocity defined over the domain  $(x, y) \in \Omega = [0, 2\pi] \times [0, 2\pi]$  with periodic boundary conditions,  $\nu = 1/Re$  is the non-dimensional viscosity equal to reciprocal of the Reynolds number and  $p$  denotes the pressure field over  $\Omega$ . We consider the flow driven by the monochromatic Kolmogorov forcing  $\mathbf{f}(\mathbf{x}) = (f_x, f_y)$  with  $f_x = \sin(k_f y)$  and  $f_y = 0$ .  $\mathbf{k}_f = (0, k_f)$  is the forcing wavenumber.

$$\begin{aligned}E(\mathbf{u}) &= \frac{1}{|\Omega|} \int_{\Omega} \frac{1}{2} |\mathbf{u}|^2 d\Omega, \\ D(\mathbf{u}) &= \frac{\nu}{|\Omega|} \int_{\Omega} |\nabla \mathbf{u}|^2 d\Omega, \\ I(\mathbf{u}) &= \frac{1}{|\Omega|} \int_{\Omega} \mathbf{u} \cdot \mathbf{f} d\Omega\end{aligned}\tag{18}$$

# Intermittent bursts of dissipation in Kolmogorov

- The Kolmogorov flow admits a laminar solution  $u_x = Re/k_f^2 \sin(k_f y)$ ,  $u_y = 0$ . For sufficiently large  $k_f$  and  $Re$ , this laminar solution is unstable, chaotic and exhibiting intermittent surges in energy input  $I$  and dissipation  $D$ .
- Study the flow under a particular set of parameters  $Re = 40$  and  $k_f = 4$  for which there is the occurrence of extreme events.



# Intermittent bursts of dissipation in Kolmogorov

Due to spatial periodicity, it is natural to examine the velocity field in Fourier space. The divergence-free velocity field  $\mathbf{u}$  admits the following Fourier series expansion:

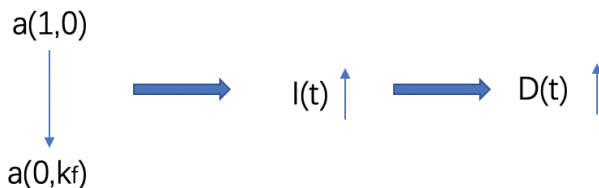
$$\mathbf{u}(\mathbf{x}, t) = \sum_{\mathbf{k}} \frac{a(\mathbf{k}, t)}{|\mathbf{k}|} \begin{pmatrix} k_2 \\ -k_1 \end{pmatrix} e^{i\mathbf{k}\cdot\mathbf{x}} \quad (19)$$

where  $\mathbf{k} = (k_1, k_2)$  is the wavenumber and  $a(\mathbf{k}, t) = \overline{-a(-\mathbf{k}, t)}$  for  $\mathbf{u}$  to be real-valued. For notation clarity, we will not explicitly write out the dependence on  $t$  from here on. Substituting [Eq \(19\)](#) into the governing equations [Eq \(17\)](#) we obtain the evolution equations for  $a$  as (more details are presented in [S1 Appendix](#))

$$\dot{a}(\mathbf{k}) = \sum_{\mathbf{p}+\mathbf{q}=\mathbf{k}} i \frac{(p_1 q_2 - p_2 q_1)(k_1 q_1 + k_2 q_2)}{|\mathbf{p}||\mathbf{q}||\mathbf{k}|} a(\mathbf{p})a(\mathbf{q}) - \nu |\mathbf{k}|^2 a(\mathbf{k}) - \frac{1}{2} i (\delta_{\mathbf{k}, \mathbf{k}_f} + \delta_{\mathbf{k}, -\mathbf{k}_f}) \quad (20)$$

# Intermittent bursts of dissipation in Kolmogorov

- In the reduced model,  
Let  $\xi_1 = a(0, k_f)$ ,  $\xi_2 = a(1, 0)$ ,  $\xi_3 = a(1, k_f)$ ,  
 $\xi_4, \xi_5, \xi_6$  are the conjugate pairs of  $\xi_1, \xi_2, \xi_3$ .
- This is because in the interest of predicting  $I$  and  $D$ , the most revealing interaction to observe is among these nodes.

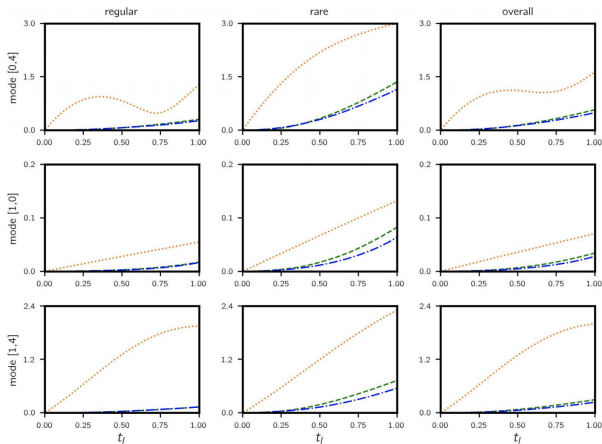


- However, only 59% energy are contained in these nodes, so a simple linear projection could lead to a bad result.

# Intermittent bursts of dissipation in Kolmogorov

- $n_{LSTM} = 70$ ,  $n_{FC} = 38$ . Number of time steps in setup stage = 100, progressively increase number of time steps in prediction stage from  $\{10, 30, 50, 100\}$ .
- 100000 trajectories, 80% training, 5% validation, 15% test.
- $n_{epochs} = 1000$ , batch size = 250,  $p_t = 60$ ,  $w_0 = 0.01$ ,  $\gamma = 0.98$ .

# Intermittent bursts of dissipation in Kolmogorov



**Fig 6. Kolmogorov flow—RMSE vs. time.** Errors are computed for  $10^4$  test trajectories (legend: fully data-driven—green dashed; data-assisted—blue dashdotted; triad—orange dotted). The RMSE in each mode is normalized by the corresponding amplitude  $E(\mathbf{k}) = \sqrt{E[|a(\mathbf{k})|^2]}$ . A test trajectory is classified as *regular* if  $|a(1,0)| > 0.4$  at  $t = 0$  and *rare* otherwise. Performance for regular, rare and all trajectories are shown in three columns. Data-assisted model has very similar errors to those of purely data-driven models for regular trajectories, but the performance is visibly improved for rare events.

# Intermittent bursts of dissipation in Kolmogorov

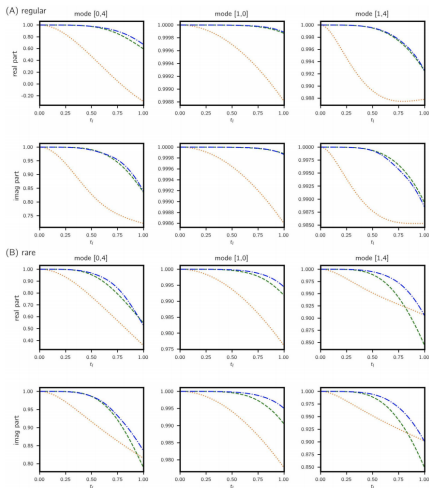
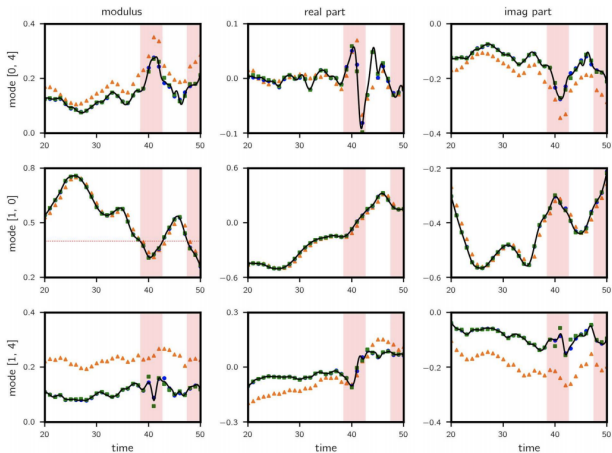


Fig 7. Kolmogorov flow: ACC vs. time. Values are computed for (A) regular and (B) rare trajectories classified from  $10^4$  test cases. Legend: fully data-driven—green dashed; data-assisted—blue dashed-dotted; trial dynamics—orange dotted. Real and imaginary parts are treated independently. Similarly to RMSE in Fig 6, improvements in predictions made by the data-assisted model are more prominent for rare events.



# Intermittent bursts of dissipation in Kolmogorov



**Fig 8. Kolmogorov flow: Predictions along a sample trajectory with lead time = 0.5.** Results for the complex modulus (left column), real part (middle column) and imaginary part (right column) of the wavenumber triad are shown. Legend: truth—black solid line; data-assisted—blue circle; triad dynamics—orange triangle; purely data-driven—green square. Rare events are recorded when  $|\text{re}(1, 0)|$  (left column, mid row) falls below 0.4 (shaded in red). Significant improvements are observed for wavenumbers (0, 4) and (1, 4).

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- Goal of the paper: Dimensionality Reduction of ODEs
- Approach: Data + ODE Prior (Physics-informed LSTM?)
- Testcases on CDV and Navier-Stokes Equations  
Also show the ability to predict rare events
- Possible future research: Singular perturbation?

# The End